#### THE JOURNAL OF CHEMICAL PHYSICS 128, 084903 (2008)

### Statistical geometry of lattice chain polymers with voids of defined shapes: Sampling with strong constraints

Ming Lin,<sup>1</sup> Rong Chen,<sup>1,2</sup> and Jie Liang<sup>2,a)</sup>

<sup>1</sup>Department of Information & Decision Science, University of Illinois at Chicago, 845 S. Morgan St., Chicago, Illinois 60607, USA <sup>2</sup>Department of Bioengineering, University of Illinois at Chicago, 851 S. Morgan St., Chicago, Illinois 60607, USA

(Received 24 October 2007; accepted 13 December 2007; published online 28 February 2008)

Proteins contain many voids, which are unfilled spaces enclosed in the interior. A few of them have shapes compatible to ligands and substrates and are important for protein functions. An important general question is how the need for maintaining functional voids is influenced by, and affects other aspects of proteins structures and properties (e.g., protein folding stability, kinetic accessibility, and evolution selection pressure). In this paper, we examine in detail the effects of maintaining voids of different shapes and sizes using two-dimensional lattice models. We study the propensity for conformations to form a void of specific shape, which is related to the entropic cost of void maintenance. We also study the location that voids of a specific shape and size tend to form, and the influence of compactness on the formation of such voids. As enumeration is infeasible for long chain polymer, a key development in this work is the design of a novel sequential Monte Carlo strategy for generating large number of sample conformations under very constraining restrictions. Our method is validated by comparing results obtained from sampling and from enumeration for short polymer chains. We succeeded in accurate estimation of entropic cost of void maintenance, with and without an increasing number of restrictive conditions, such as loops forming the wall of void with fixed length, with additionally fixed starting position in the sequence. Additionally, we have identified the key structural properties of voids that are important in determining the entropic cost of void formation. We have further developed a parametric model to predict quantitatively void entropy. Our model is highly effective, and these results indicate that voids representing functional sites can be used as an improved model for studying the evolution of protein functions and how protein function relates to protein stability. © 2008 American Institute of Physics. [DOI: 10.1063/1.2831905]

#### I. INTRODUCTION

Proteins are the working molecules of cell. Understanding how they maintain their stability and carry out their functions is a fundamental problem of molecular biology. Although it is well known that the structures of proteins are well packed,<sup>1–3</sup> there exist numerous packing defects in the form of voids buried in the interior of proteins. The size distributions of these voids are broad.<sup>4</sup> Various scaling relationships indicate that their origin may be generic steric constraints of compact chain polymers.<sup>4,5</sup> It is also well known that a few voids on a protein may play key roles in enabling protein functions,<sup>6–9</sup> for example, for substrate and ligand binding.

However, the shape space of voids of folded and unfolded proteins is not well characterized, and the energetic consequences and the kinetic effects by maintaining voids of certain shape and size are largely unknown. In this paper, we examine in detail the effects of maintaining voids of different shapes in lattice models of chain polymers. Lattice models have been widely used for studying protein folding, where the conformational space of simplified polymers can be examined in detail.<sup>10–18</sup> Despite its simplistic nature, lattice model has provided important insights about proteins, including collapse and folding transitions,<sup>16,19–23</sup> influence of packing on secondary structure and void formation,<sup>11,12,24,25</sup> the evolution of protein function,<sup>26,27</sup> nascent chain folding,<sup>18</sup> and the effects of chirality and side chains.<sup>25</sup>

In this paper, we focus on conformations that enclose voids of specific shapes. Our main objective is to study the fraction of conformations with a specific void shape among all possible conformations. This is related to the entropic cost of maintaining such a void in a polymer structure. We also study the location that voids of a specific shape and size tend to form, and the influence of compactness on the formation of such voids. The methodology we use is the sequential Monte Carlo (SMC) approach designed for sampling conformations under strong constraints, i.e., the requirement of the existence of specific types of voids. SMC is a growth-based method, in which residues are added to the chain polymer one by one until the conformation of full length is obtained. This method is first used in Ref. 28 to estimate the average extension of molecular chains. The basic goal is to obtain a set of conformational samples, along with the probabilities of generating these conformations. Compared with other sam-

128, 084903-1

<sup>&</sup>lt;sup>a)</sup>Author to whom correspondence should be addressed. Phone: (312)355-1789. FAX: (312) 996-5921. Electronic mail: jliang@uic.edu.

pling methods, such as Markov chain Monte Carlo,<sup>29–31</sup> sequential Monte Carlo can generate diverse samples and can directly estimate the number of conformations containing voids of specific shapes accurately. In this study, we develop several new strategies to improve the effectiveness of sequential Monte Carlo in generating samples under strongly constrained conditions.

Our paper is organized as follows. In Sec. II, we describe briefly the lattice model and define void and shape of voids. We then discuss the constrained sequential Monte Carlo method used in our study. Results are presented in Sec. III. The final section contains the summary and conclusion.

#### **II. METHOD**

#### A. Lattice model

In lattice models, chain polymers are self-avoiding walks (SAWs) in the square lattice space  $\mathbb{Z}^2$ . A length *n* conformation is denoted by a connected chain  $X_n = (x_1, x_2, ..., x_n)$ , where the *i*th monomer is located at the site  $x_i = (a_i, b_i)$ , where  $a_i$  and  $b_i$  are integers. The Manhattan distance between bonded monomers  $x_i$  and  $x_{i+1}$  is 1. The chain is self-avoiding:  $x_i \neq x_j$  for all  $i \neq j$ . We consider the beginning and the end of a polymer to be distinct. Only conformations that are not related by translation, rotation, and reflection are considered to be distinct. This is achieved by following the rule that a chain is always grown from the origin, the first step is always to the right, and the chain always goes up at the first time it deviates from the *x* axis. We denote the set of all length *n* SAW polymers satisfying these constraints as  $\mathcal{P}_n$ .

#### B. Voids and shape of voids

Given the conformation  $X_n \in \mathcal{P}_n$  of a chain polymer, the unoccupied sites on the square lattice are divided by the polymer into disconnected components,

$$\mathbb{Z}^2 \setminus X_n = u \cup v_1 \cdots \cup v_k,$$

where *u* is the outside component that connects to infinity, and  $v_1, \ldots, v_k$  are the *voids* that are enclosed by  $X_n$ . Here, two components are considered connected if they share any edges or vertices. By this definition, conformations (a) and (b) in Fig. 1 both have a size-2 void, but conformation (c) does not contain any void, since the internal two unoccupied points are connected to the outside through a vertex. This definition of void is arbitrary, but is consistent with the definition of contact for monomers in a chain, that is, only if two sites in a void shares an edge, they are considered to be connected.<sup>24</sup>

We are interested in the set of conformations with voids of a particular shape S. Figure 2 shows some of those shapes of sizes of 2–6. Note that those shapes are *regular* shapes, in which sites are connected by edges. In this study we do not consider shapes with sites connected by a vertex only, such as the size-2 void formed by conformation (b) in Fig. 1. Here, the voids are labeled by their shapes, where the first digit represents the number of sites the void occupies, and the second digit is the identification number of different shapes.



FIG. 1. Conformations in square lattice model. (a) This conformation also contains a size-2 void. (b) This conformation also contains a size-2 void. (c) This conformation does not contain any void.

#### C. Parameters of interest

(c)

To study the properties of conformations with specific shaped, we consider the following parameters.

#### 1. Propensity of void formation $f_1(S, n)$

Let  $\Omega_n(S)$  be the set of conformations with at least one void of a specific shape S, that is,

$$\Omega_n(S) = \{X_n | X_n$$

 $\in \mathcal{P}_n, X_n$  has at least one void of shape S}.

The fraction of conformations with void of this particular shape *S* among all possible conformations is



FIG. 2. Regular shapes of voids of different sizes. Here, the first digit represents the number of sites the void occupies, and the second digit is the identification number for different shapes. All possible shapes for voids up to size 4 are listed. Several samples for voids of sizes 5 and 6 are also listed.

$$f_1(S,n) = \frac{N(S,n)}{N_{\text{all}}(n)} = \frac{\sum_{X_n \in \Omega_n(S)} 1}{\sum_{X_n \in \mathcal{P}_n} 1}.$$
(1)

This parameter represents the propensity of void formation, i.e., the probability of forming a void of specified shape. This relates to the question whether there are preferred shapes for binding voids to occur.

# 2. Propensity of void formation with fixed loop length $f_2(I, S, n)$

The loop length of a void is defined as  $l=I_1-I_0+1$ , where  $I_0$  and  $I_1$  are the smallest and largest indices of the monomers forming the boundary of the void. Let  $\Omega_n(l,S) \subset \Omega_n(S)$  be the set of length *n* conformations with at least one void of shape *S* of loop length *l*. The fraction of conformations with void of a particular shape *S* and a particular void loop length *l* among all conformations with a void of the same shape but without the restriction of void loop length is defined as

$$f_2(l,S,n) = \frac{N(l,S,n)}{N(S,n)} = \frac{\sum_{X_n \in \Omega_n(l,S)} \xi(X_n,l,S) / K(X_n,S)}{\sum_{X_n \in \Omega_n(S)} 1}, \quad (2)$$

where  $K(X_n, S)$  is the number of shape-*S* voids in  $X_n$ ,  $\xi(X_n, l, S)$  is the number of shape-*S* voids with loop length *l* in  $X_n$ . This special treatment on N(l, S, n) is to deal with the cases when  $X_n$  has multiple voids of shape *S*. In such cases,  $X_n$  is counted once in N(S, n), and counted  $1/K(X_n, S)$  in N(l, S, n) for each combination of shape-*S* void and loop length *l*. For example, if conformation  $X_n$  has two voids of shape *S*, then  $K(X_n, S)=2$ . If both voids have loop length *l* =14, then  $\xi(X_n, l=14, S)=2$  and this conformation contributes 1 to N(l=14, S, n); if one void has loop length *l*=14 and the other void has loop length *l*=16, then this conformation contributes 1/2 to N(l=14, S, n) and 1/2 to N(l=16, S, n). Clearly, with this definition, we have

$$\sum_{l} N(l,S,n) = N(S,n)$$

The parameter  $f_2(l,s,n)$  represents the propensity of void formation with fixed loop length, i.e., the probability of forming a void of specified shape with fixed void loop

length. In protein, a related interesting question is how easier it is to form certain types of voids in shape and size with more local compared to with more global sequence fragments.

### 3. Propensity of void formation with fixed loop length and starting position $f_3(I_0, I, S, n)$

Let  $\Omega_n(I_0, l, S) \subset \Omega_n(l, S)$  be the set of length *n* conformations with at least one void of shape *S*, loop length *l*, and starting at residue position  $I_0$ . The fraction of conformations with void of a particular shape *S*, a particular loop length *l*, and a particular starting residue  $I_0$  among the conformations with a void of the same shape and the same loop length<sup>32,12</sup> is defined as

$$f_{3}(I_{0}, l, S, n) = \frac{N(I_{0}, l, S, n)}{N(l, S, n)}$$
$$= \frac{\sum_{X_{n} \in \Omega_{n}(I_{0}, l, S)} \xi(X_{n}, I_{0}, l, S) / K(X_{n}, S)}{\sum_{X_{n} \in \Omega_{n}(l, S)} \xi(X_{n}, l, S) / K(X_{n}, S)},$$
(3)

where  $\xi(X_n, I_0, l, S)$  is the number of shape-*S* voids with loop length *l* and starting residue  $I_0$  in  $X_n$ . Similarly, this definition ensures that

$$\sum_{I_0} N(I_0, l, S, n) = N(l, S, n).$$

The parameter  $f_3(I_0, I, S, n)$  represents the propensity of void formation with fixed loop length and starting position, i.e., the probability of forming a void of specified shape with fixed void loop length starting at a specific position. A related question in protein is what is the propensity of forming voids of certain shape with more local or more global sequence fragments starting at specific positions of the chain.

## 4. Propensity of void formation at specific compactness $f_4(\rho, S, n)$

The compactness of a polymer  $\rho$  is defined as  $t/t_{\max}(n)$ ,<sup>11</sup> where *t* is the number of contacts in the conformation, and  $t_{\max}(n)$  is maximum number of contacts possible for length *n* conformations. For square lattice space, we have<sup>11</sup>

$$t_{\max}(n) = \begin{cases} n - 2b, & \text{if } b^2 < n \le b(b+1) \\ n - 2b - 1, & \text{if } b(b+1) < n \le (b+1)^2 \end{cases}$$

where *b* is a positive integer. Let  $\Omega_n(\rho) \subset \mathcal{P}_n$  be the set of length *n* conformations with compactness  $\rho$  and  $\Omega_n(\rho, S) \subset \Omega_n(\rho)$  be the set of length *n* conformations with at least one void of shape *S* and compactness  $\rho$ . The fraction of conformations of a particular compactness  $\rho$  with void of a particular shape *S* among all conformations with the same compactness  $\rho$  is defined as

$$f_4(\rho, S, n) = \frac{N(\rho, S, n)}{N(\rho, n)} = \frac{\sum_{X_n \in \Omega_n(\rho, S)} 1}{\sum_{X_n \in \Omega_n(\rho)} 1}.$$
(4)

This parameter represents the propensity of void formation

Downloaded 13 Sep 2008 to 128.248.155.225. Redistribution subject to AIP license or copyright; see http://jcp.aip.org/jcp/copyright.jsp

with certain compactness, i.e., the probability of forming a void of specified shape for length n chain polymers at a fixed compactness.

# D. Estimating void parameters using sequential Monte Carlo

Exhaustive enumeration can be used to calculate the propensities defined above, but is only applicable to very short polymer chains. For longer chain, we use a modified version of the SMC method.

All the parameters described above are fractions, where the corresponding numerators and denominators N(S,n), N(l,S,n),  $N(I_0,l,S,n)$ , and  $N(\rho,S,n)$  can be written in the form of

$$\sum_{X_n \in \Omega_n(S)} h(X_n),\tag{5}$$

where  $h(\cdot)$  is a function of conformation  $X_n$ . Specifically, we have

$$\begin{split} h(X_n) &= 1 \quad \text{for } N(S,n); \\ h(X_n) &= \mathbb{I}_{\Omega_n(l,S)}(X_n) \frac{\xi(X_n,l,S)}{K(X_n,S)}, \quad \text{for } N(l,S,n); \\ h(X_n) &= \mathbb{I}_{\Omega_n(l_0,l,S)}(X_n) \frac{\xi(X_n,I_0,l,S)}{K(X_n,S)}, \quad \text{for } N(I_0,l,S,n); \\ h(X_n) &= \mathbb{I}_{\Omega_n(\rho,S)}(X_n), \quad \text{for } N(\rho,S,n), \end{split}$$

where  $\mathbb{I}_{\Omega_n}(X_n)$  is the indicator function,  $\mathbb{I}_{\Omega}(X_n)=1$  if  $X_n$  is in set  $\Omega_n$ ,  $\mathbb{I}_{\Omega}(X_n)=0$  otherwise.

Suppose we can generate random samples of conformations  $X_n^{(j)}$ , j=1, ..., m, from a trial distribution  $g(X_n)$ . Following the importance sampling principle,<sup>33,34</sup> formula (5) can be estimated as

$$\frac{1}{m}\sum_{j=1}^{m} \frac{h(X_n^{(j)})\mathbb{I}_{\Omega_n(S)}(X_n^{(j)})}{g(X_n^{(j)})} \approx \mathbb{E}_g \left[ \frac{h(X) \cdot \mathbb{I}_{\Omega_n(S)}(X_n)}{g(X)} \right]$$
$$= \sum_{X \in \Omega_n(S)} \frac{h(X)}{g(X)} \cdot g(X)$$
$$= \sum_{X_n \in \Omega_n(S)} h(X_n). \tag{6}$$

Note that to obtain an unbiased estimate, the trial distribution  $g(X_n)$  must have a support larger than  $h(X_n)\mathbb{I}_{\Omega_n(S)}(X_n)$ , that is,  $g(X_n) > 0$  must hold for all  $X_n$  in  $\Omega_n(S)$  that satisfy  $h(X_n) > 0$ . Let  $w_n^{(j)} = 1/g(X_n^{(j)})$  be the weight of sample  $X_n^{(j)}$ , then Eq. (6) can be rewritten as

$$\sum_{X_n \in \Omega_n(S)} h(X_n) = \frac{1}{m} \sum_{j=1}^m w_n^{(j)} h(X_n^{(j)}) \mathbb{I}_{\Omega_n(S)}(X_n^{(j)}).$$
(7)

The efficiency of the estimator of Eq. (6) depends on the choice of the trial distribution  $g(X_n)$  and the computational complexity for generating a sample. In general, if  $g(X_n)$  is approximately proportional to  $|h(X_n)I_{\Omega_n(S)}(X_n)|$ , with a sup-

port larger but close to  $\Omega_n(S)$ , the estimate can be reasonably accurate.<sup>34</sup>

The original Rosenbluth and Rosenbluth growth method generates samples in the space of  $\mathcal{P}_n$ .<sup>28</sup> Starting at  $x_1 = (0,0)$ , monomers are added to the chain and the associated weights are updated recursively, until the chain reaches length *n*. Modifications of the algorithm can be found in Refs. 24 and 34–36. However, the space under our consideration is a highly constrained subspace of  $\mathcal{P}_n$ . For example, for void of shape 4.1 and chain length of 50, the size of the constrained space  $\Omega_n(S)$  is less than  $2 \times 10^{-3}$  of the size of  $\mathcal{P}_n$ . With additional constraints such as fixed loop length, the space becomes more difficult. The simple growth method of<sup>28</sup> is very inefficient in generating samples for such constrained space. Below, we reformulate the sampling space and modify the growth method to overcome this difficulty.

#### 1. An equivalent representation of $\Omega_n(S)$

In order to avoid location ambiguity, the construction of  $\mathcal{P}_n$  is restricted to the set of SAW conformations starting at  $x_1 = (0,0), x_2 = (1,0)$  and going up at the first time the chain deviates from the *x* axis. Since our main interests are sampling conformations containing specific void, we adopt an equivalent representation that is more efficient for our purpose.

Specifically, let v=v(S) be a set of sites in  $\mathbb{Z}^2$ , whose union takes the shape *S*. Let  $A(v)=(a_1(v), \ldots, a_{|A(v)|}(v))$  be the set of neighboring sites of *v*, sharing either edges or vertices with *v*. We call it the *wall sites* of *v*. If a SAW completely occupies A(v) and does not intersect with *v*, then this SAW has at least one void of shape *S*. Denote the set of all such conformations as

$$G_n(v) = \{X_n | X_n \text{ is a SAW}, A(v) \subset X_n, v \cap X_n = \emptyset\}.$$

Recall that, by definition, a conformations in  $\mathcal{P}_n$  first grows to the right, and always goes up when it first deviates from the *x* axis. Note that the conformations in  $G_n(v)$  is not restricted to  $\mathcal{P}_n$  as they can start from any site on the lattice. In  $G_n(v)$ , we consider two SAWs as *equivalent* if one SAW can be transformed into the other through a combination of rotation, reflection, and position translation. Then  $G_n(v)$  consists of a number of disjoint *equivalent classes*.

It can be easily established that there is a one-to-one mapping between conformations in  $\Omega_n(S)$  and the equivalent classes of conformations in  $G_n(v)$  through transformation consisting the primitives of rotation, reflection, and translation. Each of the transformations provides such a map that the starting site  $x_1$  of  $X_n \in G_n(v)$  becomes the origin (0,0), the second site  $x_2$  becomes (1,0), and the first site that deviates from x axis is up. Hence, if  $h(\cdot)$  is a function of  $X_n$  that takes the same value for equivalent conformations, we have

$$\sum_{X_n \in \Omega_n(S)} h(X_n) = \sum_{X_n \in G_n(v)} \frac{h(X_n)}{E(X_n, S)},$$

where  $E(X_n, S)$  is the number of equivalent conformations of  $X_n$  in  $G_n(v)$ .



FIG. 3. The equivalent class of conformations. The union of the sites occupied by stars (\*) is the fixed void v. Here, polymer a, b, c, and d are different chains enclosing void v, as indicated by the different sites occupied by the starting monomer  $x_1$  and the next monomer  $x_2$ . However, the shapes of the polymers taken up by the union of the occupied sites for these chains are the same. As a consequence, these four polymers are equivalent.

The number  $E(X_n, S)$  depends on  $K(X_n, S)$ , the number of shape-S voids contained in  $X_n$  [as in Eq. (2)], and the symmetricity of the shape S. Let q(v) be the number of combination of rotation and reflection that maps v to itself. In two-dimensional lattice space, there are four possible rotations and two possible reflections around the x and y axes. Hence, q(v) can only take a value in {1,2,4,8}. For example, q(v)=8 for shape 4.2 in Fig. 2, q(v)=4 for shapes 2.1 and 3.1, q(v)=2 for shapes 4.4 and 4.5, and q(v)=1 for shape 4.3.

When  $X_n$  contains only one *S*-shaped void, the size of its equivalent class  $E(X_n, S)$  is q(v). Figure 3 shows four polymers in a equivalent class for void 2.1. When  $X_n$  contains total  $K(X_n, S)$  *S*-shaped voids, then  $E(X_n, S) = q(v)K(X_n, S)$  as each of the voids contributes q(v) number of members in the equivalent class.

To simplify our analysis, we note that  $G_n(v)$  consists of disjoint subsets,

$$G_n(v) = \bigcup_{i,k} G_n(v,i,k),$$

where

$$G_n(v,i,k) = \{X_n | X_n \in G_n(v), \quad x_k = a_i, \quad A(v) \subset \{x_1, \dots, x_k\}\}$$

If  $X_n \in G(v, i, k)$ , then the void v is completely enclosed by the prefix  $(x_1, \dots, x_k)$  of the chain, where  $x_k$  is the last monomer in the prefix and occupies the *i*th site  $a_i(v)$  of the wall sites. We have  $k \ge |A(v)|$  since some of the monomers in the prefix  $(x_1, \ldots, x_k)$  may not be on the wall of the void. The remaining chain  $(x_{k+1}, \ldots, x_n)$  does not intersect with the void space v nor with the wall sites A(v).

Using this partition, we have that for any function  $h(\cdot)$  that is constant within the equivalent classes,

$$\sum_{X_n \in \Omega_n(S)} h(X_n) = \sum_{X_n \in G_n(v)} \frac{h(X_n)}{q(v)K(X_n, S)}$$
$$= \frac{1}{q(v)} \sum_{i,k} \sum_{X_n \in G_n(v,i,k)} \frac{h(X_n)}{K(X_n, S)}.$$
(8)

In the following we develop procedures to estimate the quantity

$$\sum_{X_n \in G_n(v,i,k)} \frac{h(X_n)}{K(X_n,S)}$$

for each subset  $G_n(v,i,k)$ ,  $i=1,\ldots,|A(v)|$ ,  $k=|A(v)|,\ldots,n$ .

#### 2. Algorithmic steps

The following procedure is used to generate Monte Carlo samples in  $G_n(v, i, k)$  for all v, i, and k, which are then used to estimate the parameters listed in Sec. II C. First, we set  $x_k = a_i(v)$  as defined by  $G_n(v, i, k)$ . We then grow backwards sequentially to place  $x_{k-1}, x_{k-2}, ...$ , until we reach the



FIG. 4. The general procedure for growing chains. The union of the sites occupied by stars (\*) is the fixed void v. (a) The *k*th monomer  $y_1=x_k$  is first placed to the position  $a_i(v)$  of the wall sites of the void v. (b) We then grow backward until we reach the first monomer  $y_k=x_1$  of the chain to form void v. (c) We continue by growing forward until we reach  $x_n$ .

first monomer  $x_1$  of the chain. During this process, the wall sites A(v) become fully occupied by monomers in  $\{x_1, \ldots, x_k\}$ , and the void space v remains unoccupied. Lastly, as now that  $(x_1, \ldots, x_k)$  are placed and the constraints for void formation are satisfied, we complete the remaining conformation by sequentially placing monomers  $x_{k+1}, \ldots, x_n$ . The only constraint at this stage is that these monomers avoid the partial chain grown so far. An illustration of the procedure is shown in Fig. 4.

For ease of presentation, we rearrange the monomer labels based on the above procedure. Define  $\mathbf{y}_n = (y_1, \dots, y_n)$  as  $(x_k, \dots, x_1, x_{k+1}, \dots, x_n)$ . Formally,  $y_s = x_{k-s+1}$ , for  $s \le k$ , and  $y_s = x_s$ , for s > k. In this notation, the chain prefix of length *s* becomes  $\mathbf{y}_s = (y_1, \dots, y_s)$ .

We adopt the general framework of optimal sampling method<sup>37</sup> to generate sample conformations. Let  $m_i$  be the number of samples we retain in the *i*th iteration, and  $m_{\text{max}}$  be the maximum value of  $\{m_i\}$ . In the initial step, we set  $m_1 = 1$ ,  $y_1^{(1)} = a_i(v)$ , and  $w_1^{(1)} = m_{\text{max}}$ . For s = 2, ..., n, we perform the following procedure:

- (1) At step *s* when adding the *s*th monomer, assume there are  $m_{s-1}$  samples  $\{\mathbf{y}_{s-1}^{(j)}, j=1, \dots, m_{s-1}\}$  with weights  $w_{s-1}^{(j)}$ .
- (2) We now add the *s*th monomer to the partial chain  $\mathbf{y}_{s-1}$ . For each sample  $\mathbf{y}_{s-1}^{(j)}$ ,  $j = 1, ..., m_{s-1}$ , generate  $l_s^{(j)}$  number of new samples  $\tilde{\mathbf{y}}_s^{(l)}$  by placing  $y_s$  at each of the vacant sites neighboring  $y_{s-1}^{(j)}$ , where  $l_s^{(j)}$  is the number of vacant sites neighboring  $y_{s-1}^{(j)}$  in sample  $\mathbf{y}_{s-1}^{(j)}$ . Set weight  $\tilde{w}_{s-1}^{(l)} = w_{s-1}^{(j)}$ . Assume this step results in a total of  $L_s = \sum_j l_s^{(j)}$  samples  $(\tilde{\mathbf{y}}_s^{(l)}, \tilde{w}_s^{(l)})$ . Note that the step k+1 is slightly different. At steps 1, ..., k, we grow the chain backwards. But at step k+1, we start to grow the chain

forward. That is, we place  $y_{k+1}=x_{k+1}$ , which is connected to  $y_1=x_k$ . Hence, at the step k+1, we consider the vacant neighbor(s) of  $y_1$ , not  $y_k$ .

- the vacant neighbor(s) of y<sub>1</sub>, not y<sub>k</sub>.
  (3) Assign a priority score β<sup>(l)</sup><sub>s</sub> to each resulting partial chain ỹ<sup>(l)</sup><sub>s</sub>. The choice of the priority scores will be discussed in details in the next section.
- (4) If  $L_s < m_{\text{max}}$ , we keep all of the samples with their weights, set  $m_s = L_s$  and go to step s + 1. If  $L_s > m_{\text{max}}$ , we choose *m* distinct samples from  $\{\tilde{\mathbf{y}}_s^{(l)}, l=1, ..., L_s\}$  according to the priority scores as follows:
  - (a) Find a constant c such that  $\sum_{l=1}^{L_s} \min\{c\beta_s^{(l)}, 1\} = m_{\max}$ .
  - (b) Choose distinct integers  $J_1, J_2, ..., J_{m_{\text{max}}}$  from  $l=1, ..., L_s$ , with probability  $b_s^{(l)} \equiv \min\{c\beta_s^{(l)}, 1\}$ . This is achieved by the following steps:
  - (i) Draw a sample  $r_0$  from the uniform distribution between 0 and 1. Let  $r_i = j - r_0$  for  $j = 1, ..., m_{max}$ ;
  - (ii) For each  $j=1, ..., m_{\max}$ , choose  $J_j$  as the integer such that  $\sum_{l=1}^{J_j-1} b_s^{(l)} < r_j \leq \sum_{l=1}^{J_j} b_s^{(l)}$  holds.
- (c) Let  $\mathbf{y}_s^{(j)} = \widetilde{\mathbf{y}}_s^{(J_j)}$  and update its weight to be  $w_s^{(j)} = \widetilde{w}_s^{(J_j)} / \min\{c \boldsymbol{\beta}_s^{(J_j)}, 1\}.$

#### 3. Priority scores

The priority score guides the growth of conformations, and its design is critically important for obtaining accurate estimates. Our priority scoring function has three components addressing three important issues, namely, the support of the target distribution, the weighting scheme of samples, and the lookahead strategy.

The support of the target distribution. If a partial chain  $\tilde{\mathbf{y}}_s^{(l)}$  at step *s* is impossible to eventually grow into the constrained space  $G_n(v, i, k)$  at step *n*, it should be removed from future steps of sampling immediately at step *s*, since it is destined to be rejected. Define the support  $S_s$  of partial chains of length *s* as

$$S_s = \{\mathbf{y}_s | \text{s.t.} \exists \mathbf{y}_{s+1:n}$$
$$= (y_{s+1}, \dots, y_n) \text{ that } (\mathbf{y}_s, \mathbf{y}_{s+1:n}) \in G_n(v, i, k)\}.$$

That is,  $S_s$  contains all possible prefix chains of length *s* of desired polymers. However, it is difficult to evaluate if a partial chain is in the support. Here, we use a sequence of the support  $\psi_s$  that contains  $S_s$  but easy to work with. Specifically, let  $\psi_1 = \{a_i(v)\}$  where the chain starts according to the definition of  $G_n(v, i, k)$ . The support  $\psi_s$  is updated sequentially as follows: For each partial chain  $\mathbf{y}_{s-1} \in \psi_{s-1}$ , find all possible chains  $\mathbf{y}_s$  by adding a monomer to a vacant neighboring site which shares an edge with  $y_{s-1}$ . The new support  $\psi_s$  is the union of all such chains satisfying the following conditions:

- (i)  $y_s \cap \mathbf{y}_{s-1} = \emptyset$  (the self-avoiding constraint) and  $y_s \notin v$ , where v is the void space.
- (ii) If  $s \leq k$  and if  $A(v) \setminus \mathbf{y}_s$  is not an empty set (*i.e.*, the wall sites has not been filled by  $\mathbf{y}_s$ ), then  $A(v) \setminus \mathbf{y}_s$  must remain as a strongly connected set. Here, we define that a *strong* connection exists between two sites if they share an edge.

(iii) If  $s \le k$  and if  $A(v) \setminus \mathbf{y}_s \ne \emptyset$ , then the site  $y_s$  must satisfy

 $k-s \ge d(y_s, A(v) \setminus \mathbf{y}_s) + |A(v) \setminus \mathbf{y}_s| - 1,$ 

where  $d(y_s, A(v) \setminus \mathbf{y}_s)$  is the minimum Manhattan distance between  $y_s$  and the unoccupied wall sites,  $|A(v) \setminus \mathbf{y}_s|$  is the number of unoccupied wall sites of  $A(v) \setminus \mathbf{y}_s$ .

Condition (ii) reflects the property that both the filled and unfilled sites on the wall of the void must remain strongly connected at any time of the growth. Otherwise, the unfilled wall sites A(v) has multiple not strongly connected components. In such cases, the self-avoiding property must be violated in order to fill all of them by  $\mathbf{y}_n \in G_n(v, i, k)$ . This is the consequence of the Jordan curve theorem in plane.<sup>38</sup>

Condition (iii) is to ensure that the remaining chain length is sufficient to fill all wall sites, i.e.,  $A(v) \setminus y_s$  must be filled by  $(y_{s+1}, \ldots, y_k)$ , which is a length k-s chain connected to  $y_s$ .

The priority scores without lookahead. In the optimal sampling method framework,<sup>37</sup> the priority score serves both as the propagation trial distribution as well as the resampling priority score. Under the importance sampling principle,<sup>34</sup> the ideal trial distribution should be proportional to  $|h(x)\pi(x)|$ , where  $\pi(x)$  is the target distribution. In our case, it translates to  $\widetilde{w}_t^{(l)}h(\widetilde{\mathbf{y}}_t^{(l)})\mathbb{I}_{\psi_t}(\widetilde{\mathbf{y}}_t^{(l)})$ . Here,  $h(\widetilde{\mathbf{y}}_t^{(l)})$  is the value of function  $h(\cdot)$  applied to partial chain  $\widetilde{\mathbf{y}}_t^{(l)}$ , which is always non-negative.

For s > k, we simply set equally priority scores  $\beta_s^{(l)} \equiv 1.0$  to all partial chain samples  $\tilde{\mathbf{y}}_s^{(l)}$ , since this stage is relatively easy.

For the more difficult part of the growth  $s \le k$  where the major constraints lie, we need to guide samples to grow into the support region  $\psi_s$  in order to reduce the sample rejection rate. Following Zhang and Liu,<sup>36</sup> we use the priority score to achieve this. Taking condition (iii) when updating the support into consideration, we define

$$U_{s}(\widetilde{\mathbf{y}}_{s}^{(l)}, v) = k - s + 2 - d(\widetilde{y}_{s}^{(l)}, A(v) \setminus \widetilde{\mathbf{y}}_{s}^{(l)}) - |A(v) \setminus \widetilde{\mathbf{y}}_{t}^{(l)}|$$

if  $|A(v) \setminus \tilde{\mathbf{y}}_t^{(l)}| > 0$ ,  $U_s(\tilde{\mathbf{y}}_s^{(l)}, v) = 0$  if  $|A(v) \setminus \tilde{\mathbf{y}}_t^{(l)}| = 0$ . It evaluates how much freedom and flexibility the remaining chain possesses. When  $|A(v) \setminus \tilde{\mathbf{y}}_t^{(l)}| > 0$ , there are still some vacant sites on the void wall needs to be occupied. In this case, if  $U_s(\tilde{\mathbf{y}}_s^{(l)}, v) \leq 0$ ,  $\tilde{\mathbf{y}}_s^{(l)}$  is not in the support  $\psi_s$ , as it violates condition (iii), we reject this sample. The larger  $U_s(\tilde{\mathbf{y}}_s^{(l)}, v)$  is, the less constrained the remaining chain is.

Combining the value of the function  $h(\tilde{\mathbf{y}}_t^{(l)})$  to be evaluated, and  $U_s(\tilde{\mathbf{y}}_s^{(l)}, v)$  reflecting the desired flexibility of the remaining chain, we design our priority score for  $s \leq k$  as

$$\boldsymbol{\beta}_{s}^{(l)} = \widetilde{w}_{s}^{(l)} h(\widetilde{\mathbf{y}}_{s}^{(l)}) \mathbb{I}_{\psi_{s}}(\widetilde{\mathbf{y}}_{s}^{(l)}) \exp\{-U_{s}^{-1/2}(\widetilde{\mathbf{y}}_{s}^{(l)}, \upsilon)/T_{s}\},\$$

where  $T_s$  is a temperaturelike variable. The choice of values for  $T_s$  is important. In general, the constraint of forming void is not of serious concerns at the beginning, so we can use high values of  $T_s$  to enhance diversity in sampling. As the chain grows, the concern of meeting the constraints become stronger, since there are less freedom for the remaining chains to grow. Hence, we gradually reduced the  $T_s$ , as in simulated annealing algorithms. In this study, we use  $T_s = \sqrt{k-s+16}$  for  $s=1, \ldots, k-1$ .

*Priority score with lookahead.* An often used strategy to improve performance of SMC is lookahead.<sup>36,39,40</sup> Lookahead enables us to use information from possible future steps to construct priority scores, resulting smaller rejection rate of the samples. In addition, it reduces the variance of samples for estimation and hence improves sample efficiency.<sup>41</sup>

For a  $\delta$ -step lookahead, the priority score at time *t* is determined by exploring all possible combinations of  $\delta$ -step growth from the current sample  $\mathbf{y}_s$ . Specifically, the priority scores we use are

$$\beta_{s}^{(l)} = \widetilde{w}_{s}^{(l)} \sum_{y_{s+1}, \dots, y_{s+\delta}} h(\widetilde{\mathbf{y}}_{s+\delta}^{(l)}) \mathbb{I}_{\psi_{s+\delta}}(\widetilde{\mathbf{y}}_{s+\delta}^{(l)}) \times \exp\left\{-\frac{U_{s}^{-1/2}(\widetilde{\mathbf{y}}_{s+\delta}^{(l)}, v)}{T_{s}}\right\},$$

where  $\tilde{\mathbf{y}}_{s+\delta}$  denotes  $(\tilde{\mathbf{y}}_s, y_{s+1}, \dots, y_{s+\delta})$ .

Note that as lookahead step  $\delta$  increases, the effectiveness increases at the cost of exponentially growing computational complexity. Hence, the choice of  $\delta$  is a tradeoff between estimate efficiency and complexity. In this study we use  $\delta = 1$ .

#### 4. Estimation

In our framework, it is possible to estimate the parameters described in Sec. II C for polymer chains of different lengths up to n when generating conformation samples of length n.

Specifically, when generating conformation samples for  $G_n(i,k)$ , at step s=k, k+1, ..., n, the generated partial conformations are  $\tilde{\mathbf{y}}_s^{(l)} = (x_1^{(l)}, ..., x_k^{(l)}, ..., x_s^{(l)})$ , which are properly weighted chain polymers of length *s*. Hence,  $\sum_{\mathbf{x}_{n^*} \in G_{n^*}(v,i,k)} h(\mathbf{x}_{n^*})/K(\mathbf{x}_{n^*}, S)$ ,  $n^*=k, k+1, ..., n$ , can be estimated by the following estimator:

$$\hat{h}(\mathbf{x}_{n^*}; i, k) = \frac{1}{m_{\max}} \sum_{l=1}^{L_s} \widetilde{w}_s^{(l)} \frac{\mathbb{I}_{G_{n^*}(v, i, k)}(\widetilde{\mathbf{y}}_s^{(l)}) h(\widetilde{\mathbf{y}}_s^{(l)})}{K(\widetilde{\mathbf{y}}_s^{(l)}, S)}$$

at step  $s=n^*$ . Here, the estimation is made after step (2) of the algorithmic steps in the previous subsection.

After generating samples for  $G_n(i,k)$  of all possible i,k, for any  $n^* \leq n$ , we can estimate  $\sum_{X_{n^*} \in \Omega_{n^*}(S)} h(X_{n^*})$  by

$$\sum_{X_{n^*}\in\Omega_{n^*}(S)}h(X_{n^*})\approx\frac{1}{q(v)}\sum_{i,k}\hat{h}(\mathbf{x}_{n^*};i,k),$$

according to Eq. (8).

#### **III. RESULTS**

In this section, we present the results of estimation of the parameters described in Sec. II C. We also develop a parametric model relating to void and chain properties for interpreting the estimated results and for prediction of propensity of forming void of specific shape.



FIG. 5. Estimating void propensity values. (a) Estimated propensity values and true propensity values of forming size-4 voids of different specific shapes for conformations of lengths of 14–24. They superimpose very well. (b) Estimated propensity values of forming size-4 voids of different specific shapes for conformations of lengths of 15–50.

#### A. Propensity of void formation

For propensity of void formation  $f_1(S,n)$  defined in Eq. (1), we first examine size-4 voids. There are five different shapes for size-4 regular voids (Fig. 2). To validate our procedure, the estimated propensity of void formation is compared with the true values obtained by exhaustive enumeration, for chains of length of 14–24. Figure 5(a) shows the results for voids shapes 4.3 and 4.4. The estimated values are indistinguishable from the true values. These results suggest that our sampling method works well and can provide accurate estimations.

The results for longer chains of length of 15–50 using the SMC procedure are presented in Fig. 5(b), where propensity of void formation  $f_1(S,n)$  for void shapes 4.1–4.5 are shown. It is clear that voids of different shapes have significant difference in propensity of formation. This raises the question whether voids and binding sites in proteins are simi-



FIG. 6. Propensity values of forming voids (size=2-5, a-d) of different specific shapes for conformations of lengths of 25–50. These are used to develop a regression model. Dashed line: results obtained by estimation using sequential Monte Carlo. Solid line: fitted results from the regression models.

larly biased, and whether the distribution of voids of different shapes can be partly explained by these intrinsic propensities analogous to what is observed here on lattice models.

#### 1. Predictive models

To better understand our estimation results and to infer general principles, we develop a predictive model for  $f_1(S,n)$ using the following parametric form:

$$\hat{f}_1(S,n) = \frac{1}{q(v)} c_1 c_2^{-|A(v)|} (n - |A(v)| + 1)^{c_3} [1 - c_4(|e(v)| - 4)],$$
(9)

where q(v) represents the degeneracy of the void shape as we discussed in Sec. II D 1. We consider three factors other than q(v) in our model: the the wall size |A(v)|, the chain length *n*, and the number of outer corners of void, |e(v)|. Here, the outer corners, e(v), are defined as the sites on void wall that connect to the void through a single vertex only. The values of q(v), |A(v)|, and |e(v)| for different void shapes are summarized in Table I.

In this model,  $c_1$ ,  $c_2$ ,  $c_3$ , and  $c_4$  are positive constants. As the wall size |A(v)| increases, it is expected that the propensity of forming voids of the specific shape decreases exponentially. This is reflected by the term containing  $c_2^{-[A(v)]}$ . When the chain length *n* increases, it is expected that the propensity of forming voids of the specific shape increases by some power. This is captured by the term of

TABLE I. Geometric features of voids determining the fractions of chain polymers containing such voids. q(v) is related to the symmetry of the void v, |A(v)| is wall size of the void v, and |e(v)| is the number of outer corners of void v.

Void type	2.1	3.1	3.2	4.1	4.2	4.3	4.4	4.5	5.1	6.1	6.2	6.3
q(v)	4	4	2	4	8	1	2	2	4	4	1	2
A(v)	10	12	12	14	12	14	14	14	16	18	18	18
e(v)	4	4	5	4	4	5	6	6	4	4	5	6

Downloaded 13 Sep 2008 to 128.248.155.225. Redistribution subject to AIP license or copyright; see http://jcp.aip.org/jcp/copyright.jsp



FIG. 7. Estimated and predicted propensity values of forming size-6 voids of different specific shapes for conformations of lengths of 25–50. Dashed line: SMC results. Solid line: predicted results using the regression model (9).

 $(n-|A(v)|+1)^{c_3}$ . We also find that the number of outer corners, |e(v)|, is an important determinant of propensity of void formation. For void shapes with more outer corners, chain polymers enclosing such voids have more concave turns on the wall. This makes it more difficult for a self-avoiding chain to enclose the void. The negative term of |e(v)| in Eq. (9) models this effect.

We estimate the coefficients in model (9) using the estimated  $f_1(v_n)$  from SMC for voids of sizes 2–5 and chain length from 25 to 50. Taking log transformation and using nonlinear regression, we found that  $\hat{c}_1$ =47.46,  $\hat{c}_2$ =2.28,  $\hat{c}_3$ =0.76, and  $\hat{c}_4$ =0.21.



FIG. 8. Estimated propensity values of forming size-4 voids of different specific shapes with different specified loop length for conformations of length n=50.



FIG. 9. Conformation with odd loop length. This conformation encloses an void of shape 4.1 with loop length of 17.

The propensity values estimated from SMC and the fitted results of  $\hat{f}_1(S,n)$  using model (9) are plotted in Fig. 6. It can be seen that the parametric model fits the data very well. Using the above estimated parameters obtained from the training data, we develop a predictive model for the propensity for void shapes 6.1, 6.2, and 6.3, which are not used in deriving the regression model. The predictions are again compared with those estimated by SMC (Fig. 7). The models works well, although it consistently underestimates by a small amount for void shape 6.3.

#### B. Propensity of void formation with fixed loop length

Now we consider the propensity of void formation with fixed loop length  $f_2(S,n)$  defined in Eq. (2). We plot estimated  $f_2(l, S, n=50)$  for different specified loop lengths *l* and shapes *S* in Fig. 8. Although voids with odd loop length do exist, we can see that it is much easier to form void, with even loop length. This is because the number of wall sites, |A(v)|, is always an even number on lattice. To form a void *v* with odd loop length, the first monomer and the last monomer of the polymer on the void wall A(v) cannot be adjacent, which results in a more complicated shape. A conformation enclosing a void of shape 4.1 with loop length of 17 is given in Fig. 9. On average, void shapes 4.1 and 4.3 have larger loop sizes than void shapes 4.4 and 4.5, because they have fewer corners. These results suggest that voids of different shapes have different propensity at specific loop lengths.

#### C. End effect for void formation

For propensity of void formation with fixed loop length and specified starting position  $f_3(I_0, l, S, n)$ , as defined in



FIG. 10. Estimated propensity values of forming size-4 voids of different specific shapes with fixed loop length l=14 and different specified starting position for conformations of length n=50.



FIG. 11. Estimated propensity values of forming size-4 voids of different specific shapes with certain compactness for conformations of length from n=30 to 50. [(a)–(d)] void 4.1, void 4.3, void 4.4, and void 4.5.

Eq. (3), we plot estimated  $f_3(I_0, l=14, S, n=50)$  for voids of shapes *S* with a loop length of 14 in chain polymers of length 50 with different starting positions  $I_0$  in Fig. 10. We find that the propensities  $f_3(I_0, l, S, n)$  at  $I_0=1$  and  $I_0=2$  are very different, indicating strong end effect for void formation. That is, void is much easier to form at the end of the conformation. This is likely due to the *tail effect*. Void at the end of the chain only need to have one tail, but has two tails if it is in the middle of the conformation. It is difficult to constrain the tails to satisfy the multiple restrictions for forming void of certain shapes.

# D. Propensity of void formation at different compactness

Figure 11 shows estimated propensity values of void formation at different compactness  $f_4(\rho, S, n)$  defined in Eq. (4) for chain length from n=30 to 50. Conformations with size-4 voids are dominated by those at compactness around 0.3–0.7. If we normalize  $f_4(\rho, S, n=50)$ , that is, we define

$$\overline{f}_4(\rho, S, n) = \frac{f_4(\rho, S, n)}{\int f_4(\rho, S, n) d\rho}$$

where  $f_4(\rho, S, n)$  can be considered as a distribution of  $\rho$ . We plot the 0.25 quantile, median value, and 0.75 quantile of distribution  $\overline{f}_4(\rho, S, n)$  for different chain length *n* and fixed shape *S* in Fig. 12. We can see these values slightly increase as *n* increases from 30 to 50. This indicates that the preferred

compactness range of forming these size-4 voids shifts slightly to more compact regions as chain length increases. We also compare the propensity values of forming voids of all size-2 regular shapes (2.1), voids of all size-3 regular shapes (3.1, 3.2), and voids of all size-4 regular shape (4.1, 4.2, 4.3, 4.4, 4.5) for chains of length of 50 at different compactnesses (Fig. 13). The results show that smaller voids are easier to form as compactness increases. Our results from lattice model suggests that there might be a preferred size for void formation in proteins, which are all within a specific narrow range of compactness.<sup>3</sup>

#### **IV. SUMMARY AND CONCLUSION**

Protein molecules contain many voids buried in the interior of proteins, with broad distribution.<sup>4</sup> Although most voids are likely to originate from generic steric constraints of compact chain polymers,<sup>4,5</sup> some voids are the functional regions for many proteins, such as enzymes, where substrates and ligands bind and biochemical reactions occur.<sup>6,7</sup>

An important general question is how the need for maintaining functional voids, which have to be of specific shape, is influenced by, and affects other aspects of proteins structures and properties: e.g., protein folding stability, kinetic accessibility, and evolutionary selection pressure. These are broad and complex issues that require detailed studies.

In this work, we study the effects of maintaining voids of defined shape using lattice model. Because the conformational space of simplified polymers can be examined in de-



FIG. 12. Estimated quantiles (0.25, 0.5, and 0.75) of distribution  $\overline{f}_4(\rho, S, n)$  for different chain length *n* and void shape *S*. [(a)–(d)] void 4.1, void 4.3, void 4.4, and void 4.5.

tail, lattice models have been widely used in protein studies and have lead to important insight about protein folding. The focus of our study is to generate large number of sample conformations under very constraining restrictions to study general properties of voids and their shapes. We use sequential Monte Carlo method and have developed an efficient growth method to generate conformation samples in highly constrained space.

We show that our approach is effective in estimating entropy of void maintenance, with and without an increasing number of restrictive conditions, such as loops forming the wall of void with fixed length, with additionally fixed starting position in the sequence. Our results also lead to a num-



FIG. 13. Estimated propensity values of forming voids of all size-2 regular shapes (2.1), voids of all size-3 regular shapes (3.1, 3.2), and voids of all size-4 regular shape (4.1, 4.2, 4.3, 4.4, 4.5) for chain length of 50 and different compactness.

ber of observations, including that polymers of certain compactness range favors the formation of voids of specific size, and that voids are far easier to form around the end of the polymer. A finding is that voids tend to form at the chain ends. This raises the interesting question whether voids and pockets tend to form at either the N-terminal or the C-terminal end in real proteins. A detailed analysis of voids and pockets in real proteins will be necessary for answering this question. In addition, we have developed a parametric model for explaining the propensity of forming voids of particular shapes, or equivalently, the entropic cost of maintaining such voids. Our model is highly effective in predicting the propensity of void formation for different shapes. Such lattice model of voids representing functional sites can be used as improved model for studying the evolution of protein functions,<sup>26</sup> and how it relates to protein stability.<sup>2</sup>

Although in this study we treat the occurrence of all conformations equally likely, our approach can be applied to models with more realistic energy functions in a straightforward manner. The approach for sampling strongly constrained conformations we developed in this study will be generally applicable for studying real proteins in threedimensional space.

- <sup>1</sup>F. M. Richards, Annu. Rev. Biophys. Bioeng. 6, 151 (1977).
- <sup>2</sup>C. Chothia, Nature (London) **254**, 304 (1975).
- <sup>3</sup>F. M. Richards and W. A. Lim, Q. Rev. Biophys. 26, 423 (1994).
- <sup>4</sup>J. Liang and K. A. Dill, Biophys. J. **81**, 751 (2001).
- <sup>5</sup>J. Zhang, R. Chen, C. Tang, and J. Liang, J. Chem. Phys. **118**, 6102 (2003).
- <sup>6</sup>R. A. Laskowski, N. M. Luscombe, M. B. Swindells, and J. M. Thornton,

- <sup>7</sup>J. Liang, H. Edelsbrunner, and C. Woodward, Protein Sci. **7**, 1884 (1998).
- <sup>8</sup>T. A. Binkowski, L. Adamian, and J. Liang, J. Mol. Biol. **332**, 505 (2003).
- <sup>9</sup>Y. Tseng and J. Liang, Mol. Biol. Evol. **23**, 421 (2006).
- <sup>10</sup>K. F. Lau and K. A. Dill, Macromolecules **93**, 6737 (1989).
- <sup>11</sup>H. S. Chan and K. A. Dill, Macromolecules 22, 4559 (1989).
- <sup>12</sup>H. S. Chan and K. A. Dill, J. Chem. Phys. **92**, 3118 (1990).
- <sup>13</sup>E. Shakhnovich and A. Gutin, J. Chem. Phys. **93**, 5967 (1990).
- <sup>14</sup>C. J. Camacho and D. Thirumalai, Proc. Natl. Acad. Sci. U.S.A. **90**, 6369 (1993).
- <sup>15</sup> V. S. Pande, C. Joerg, A. Y. Grosberg, and T. Tanaka, J. Phys. A **27**, 6231 (1994).
- <sup>16</sup>N. D. Socci and J. N. Onuchic, J. Chem. Phys. **101**, 1519 (1994).
- <sup>17</sup> K. Dill, S. Bromberg, K. Yue, K. Fiebig, D. Yee, P. Thomas, and H. Chan, Protein Sci. 4, 561 (1995).
- <sup>18</sup> H. Lu and J. Liang, "A model study of protein nascent chain and cotranslational folding using hydrophobic-polar residues," Proteins (to be published).
- <sup>19</sup>A. Šali, E. I. Shakhnovich, and M. Karplus, Nature (London) **369**, 248 (1994).
- <sup>20</sup> I. Shrivastava, S. Vishveshwara, M. Cieplak, A. Maritan, and J. R. Banavar, Proc. Natl. Acad. Sci. U.S.A. **92**, 9206 (1995).
- <sup>21</sup>D. K. Klimov and D. Thirumalai, Phys. Rev. Lett. 76, 4070 (1996).
- <sup>22</sup> R. Mélin, H. Li, N. Wingreen, and C. Tang, J. Chem. Phys. **110**, 1252 (1999).
- <sup>23</sup>S. Kachalo, H. Lu, and J. Liang, Phys. Rev. Lett. **96**, 058106 (2006).

J. Chem. Phys. 128, 084903 (2008)

- <sup>24</sup>J. Liang, J. Zhang, and R. Chen, J. Chem. Phys. **117**, 3511 (2002).
- <sup>25</sup> J. Zhang, Y. Chen, R. Chen, and J. Liang, J. Chem. Phys. **121**, 592 (2004).
- <sup>26</sup> P. D. Williams, D. D. Pollock, and R. Goldstein, J. Mol. Graphics Modell. 19, 150 (2001).
- <sup>27</sup> J. D. Bloom, C. O. Wilke, F. H. Arnold, and C. Adami, Biophys. J. 86, 2758 (2004).
- <sup>28</sup> M. N. Rosenbluth and A. W. Rosenbluth, J. Chem. Phys. **23**, 356 (1955).
- <sup>29</sup> W. R. Gilks, S. Richardson, and D. J. Spiegelhalter, *Markov Chain Monte Carlo in Practice* (Chapman and Hall, London, 1996).
- <sup>30</sup>T. Hamelryck, J. Kent, and A. Krogh, PLOS Comput. Biol. **2**, 1121 (2006).
- <sup>31</sup> Y. Iba, G. Chikenji, and M. Kikuchi, J. Phys. Soc. Jpn. **67**, 3327 (1998).
- <sup>32</sup>H. S. Chan and K. A. Dill, J. Chem. Phys. **90**, 492 (1989).
- <sup>33</sup>A. Marshall, in *Symposium on Monte Carlo Methods*, edited by M. Meyer (Wiley, New York, 1956), pp. 123–140.
- <sup>34</sup> J. S. Liu, *Monte Carlo Strategies in Scientific Computing* (Springer, New York, 2001).
- <sup>35</sup> P. Grassberger, Phys. Rev. E 56, 3682 (1997).
- <sup>36</sup>J. L. Zhang and J. S. Liu, J. Chem. Phys. **117**, 3492 (2002).
- <sup>37</sup> P. Fearnhead and P. Clifford, J. R. Stat. Soc. Ser. B (Stat. Methodol.) 65, 887 (2003).
- <sup>38</sup>A. Hatcher, *Algebraic Topology* (Cambridge University Press, Cambridge, England, 2002).
- <sup>39</sup>H. Meirovitch, J. Phys. A **15**, L735 (1982).
- <sup>40</sup>X. Wang, R. Chen, and D. Guo, IEEE Trans. Signal Process. **50**, 241 (2002).
- <sup>41</sup>A. Kong, J. Liu, and W. Wong, J. Am. Stat. Assoc. **89**, 278 (1994).

Protein Sci. 5, 2438 (1996).